

Vibrationally Induced Two-Level Systems in Single-Molecule Junctions

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Single-molecule junctions are found to show anomalous spikes in dI/dV spectra. The position in energy of the spikes are related to local vibration mode energies. A model of vibrationally induced two-level systems reproduces the data very well. This mechanism is expected to be quite general for single-molecule junctions. It acts as an intrinsic amplification mechanism for local vibration mode features and may be exploited as a new spectroscopic tool.

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A single atom or molecule with an almost transparent single conductance channel leading to a conductance near the conductance quantum $2e^2/h$ ($= 1 G_0$) can be contacted to leads. Conduction electrons can pass through such junction ballistically for low bias voltages since the mean free path of the electrons is much larger than the size of the contact. However the contact is not entirely ballistic in the sense that once the excess energy of the conduction electrons becomes equal or larger than the energy of a local mode of the contact, the electrons can scatter inelastically by exciting a local mode. This results in the case of a perfectly transmitting single channel contact to a small decrease in the conductance, since the forward travelling electrons are backscattered due to the energy loss in the inelastic scattering process. Differential conductance (dI/dV) measurements have identified vibration modes of single molecules in an atomic contact [1]. This technique, also called Point Contact Spectroscopy (PCS) is analogous to inelastic electron tunnelling spectroscopy (IETS) for single molecules [2, 3], with the difference that the conductance in the latter case increases due to the opening of an additional conductance channel.

In this letter we present the observation of anomalous spikes, rather than steps, in dI/dV measurements on various single molecule contacts. We present a model that involves two-level systems, which describes our data very well. It may be used as a new spectroscopic tool for identifying molecular vibration modes in single molecule junctions.

We create atomic contacts using a mechanically controlled break junction (MCBJ) setup in cryogenic vacuum at 4.2 K (see ref. [4] for a detailed description). Break junctions for the metals Au, Ag, Pt and Ni have been investigated with the molecules H_2 , D_2 , O_2 , C_2H_2 , CO , H_2O and benzene. In most of these cases regular vibration mode spectra displaying a step down in conductance have been observed, but for all systems anomalous spectral features as displayed in Fig. 1 were also found.

In order to admit these molecules to the metal atomic contacts at 4.2 K, the insert is equipped with a capillary that has a heating wire running all along its interior to prevent premature condensation of the gasses. The amount of gas admitted is of order $10 \mu\text{mol}$. Previous

measurements have clearly demonstrated that it is possible to capture a single molecule in the atomic junction and measure its vibration modes [1, 5].

Figure 1(a) shows a dI/dV spectrum of a single CO molecule contacted by Pt leads. The observation of spikes that dominate the spectrum is surprising. Pt-CO-Pt contacts at a conductance of $G \simeq 1G_0$, as for Pt-H-H-Pt contacts [1], typically display step-like decreases in dI/dV when conduction electrons are backscattered due to energy loss in exciting a local vibration mode [7]. Figure 2(c) shows an example of a regular dI/dV spectrum for a Pt-D-D-Pt contact. The I-V curve of Fig. 1(b) shows clear fluctuations in the voltage window in which the peaks in dI/dV appear. At high voltages the fluctuations have disappeared and the junction displays a slightly higher conductance.

We have evidence that the anomalies are related to vibration modes of the single molecule. Firstly, for the Pt-H₂ system the spikes in dI/dV typically appear at the same energies as the transverse vibration mode energies of a hydrogen molecule [5]. Figures 2(a) to (c) show a sequence of spectra taken when a Pt-D-D-Pt contact was slightly stretched over 0.2 Å. The small bumps on the vibration mode shoulders in Fig. 2(a) at ± 48 mV evolve into full peaks in Fig. 2(b) and subsequently into a regular vibration mode spectrum for a deuterium molecule in between Pt contacts in Fig. 2(c). Additionally we have introduced hydrogen to atomic gold contacts and

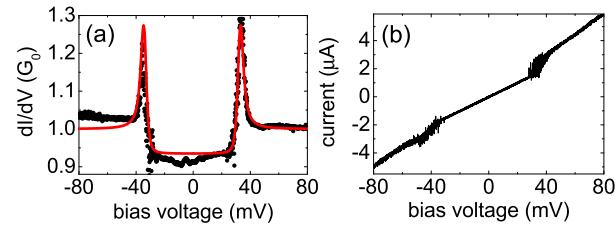


FIG. 1: (Color online) (a) dI/dV spectrum on a CO molecule bridging a Pt contact displaying symmetric positive spikes. The red curve shows a fit by the model discussed in the text [6]; (b) I-V measurement on a CO molecule bridging a Pt contact

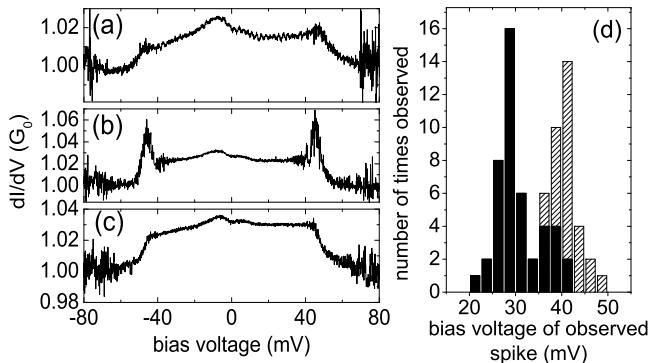


FIG. 2: (a) to (c) display the gradual change of dI/dV spectra of a Pt-D-D-Pt contact; (b) and (c) show spectra taken for the same junction as (a) after stretching by 0.1 \AA and 0.2 \AA , respectively. Histogram (d) of peak positions in spectra of junctions for Au-H₂ (hatched bars) and Au-D₂ (filled bars). All data were obtained at $T = 4.2\text{K}$.

subsequently measured dI/dV . The spectra show similar anomalous spikes as for Pt-H₂. The positions of the spikes observed in dI/dV spectra for Au-H₂ and Au-D₂ are shown as histograms in Fig. 2(d). It can be clearly seen that the histogram maximum for Au-D₂ at 29 meV, has shifted by the square root of the mass ratio from 42 meV for Au-H₂. From the evidence in Fig. 2 it can be concluded that the observed singularities are associated with the molecular vibration mode energies in the contact.

Spikes in dI/dV point to abrupt changes in the conductance of the contacts in small voltage windows. Since similar spikes have been observed for different molecules this suggests that it involves a general feature of single molecule junctions. The ‘noise’ visible in Fig. 1(b) is a result of two-level fluctuations, that are excited at high bias. However, they are apparently not excited directly in view of the connection with vibration mode energies shown in Fig. 2. Reference [8] has shown that by exciting internal vibrations of a pyrrolidine molecule on a metal surface, the molecule can flip to a different conformation. The different conformations have different conductances leading to a dip in dI/dV . This forms a basis for our proposed explanation.

We attribute the abrupt changes in contact conductance to abrupt switching between two slightly different local geometrical configurations that have different conductances. Current induced two-level systems in mesoscopic contacts have been studied in detail in the past (see e.g. [9, 10, 11]). In those cases spikes are only observed in the second derivative d^2I/dV^2 , in contrast to spikes in dI/dV in our measurements. Since the anomalous spikes appear around the vibration mode energies, we assume that the switching is induced by intramolecular vibrations. STM studies on single molecules on metal surfaces have clearly shown, e.g., that by excit-

ing an internal vibration mode of a molecule the lateral hopping rate of a CO molecule [12], the rotation of a O₂ molecule [13] or of a C₂H₂ molecule [14] can be increased by two orders of magnitude. In our case the molecules are chemically linked by two leads preventing the molecules of jumping away and the current is much higher, leading to orders of magnitude higher rates.

We derive a simple model of a vibrationally induced two-level system (VITLS) that reproduces the shape of the anomalous spikes in our dI/dV spectra very well. The model is schematically depicted in Fig. 3. The molecule experiences a double well potential landscape with ground states Ψ_0 and Ψ_1 in the two energy minima. The two energy minima are separated by a large potential barrier that prevents tunneling between the two ground states. The molecule in the contact can be vibrationally excited by conduction electrons to the state Ψ^e . In order to displace the molecule, the barrier for the reaction coordinate has to be overcome. The direction of the internal vibration needs to be in the direction of the reaction coordinate (i.e. the movement) of the molecule [15]. In the case of the Pt-D₂ junctions the spikes have only been observed at energies corresponding to vibrations perpendicular to the current direction (i.e. transverse modes). The longitudinal vibration mode at about 90 meV [5] never showed anomalous spikes. In our quasi one-dimensional junctions the transverse modes are the ones that have the proper reaction coordinates to induce changes in the geometry of the junctions.

The expression for the current passing through the junction can be approximately written as:

$$I(V) = \left(n_0\sigma_0 + n_1\sigma_1 + n_e \frac{(\sigma_0 + \sigma_1)}{2} \right) V \quad (1)$$

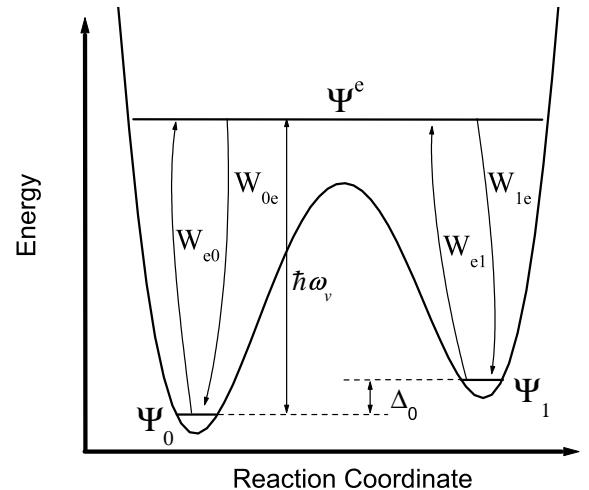


FIG. 3: Energy landscape of a molecule in an atomic contact. The molecule can be vibrationally excited and relax into the other energy minimum.

We assume that the conductance σ_0 of well 0 and σ_1 of well 1 are different. A step in $I(V)$ and thus a peak in dI/dV can be obtained once Ψ_1 becomes populated. The total current is determined by the sum of the products of the occupation numbers with their respective conductances. For simplicity we take the average of σ_0 and σ_1 to be the conductance of the excited state. In order to determine the shape of the dI/dV spectrum, the dependence of n_0 , n_1 and n_e on V has to be determined. Let us therefore consider dynamic equilibrium of the occupation numbers of all levels involved in the system as displayed in Fig. 3 in terms of the transition rates,

$$\begin{aligned} \frac{dn_0}{dt} &= -W_{e0} + W_{0e} = 0 \\ \frac{dn_1}{dt} &= -W_{e1} + W_{1e} = 0 \\ \frac{dn_e}{dt} &= -W_{0e} - W_{1e} + W_{e0} + W_{e1} = 0 \\ n_0 + n_1 + n_e &= 1. \end{aligned} \quad (2)$$

Here W_{fi} represent the transition rates from states i to states f . The energy difference between the two ground states (Fig. 3) we take to be Δ_0 . The vibration mode energy $\hbar\omega_v$ is much larger than Δ_0 and $k_B T \leq \Delta_0$ to keep the system well defined in the ground state for low bias voltages. The transition rates W can now be determined. Consider as an example the rate W_{e0} :

$$W_{e0} = n_0 \int_{-\infty}^{+\infty} \gamma f(E, eV) [1 - f(E - \hbar\omega_v, eV)] dE \quad (3)$$

We assumed here that the electron density of states is flat over the energy window under consideration. Where γ is the electron-phonon coupling and $f(E, eV)$ are the non-equilibrium electron distribution functions. All other rates can be expressed in a similar way. After considering the low temperature limit, $\frac{\hbar\omega_v}{k_B T}, \frac{(\hbar\omega_v - \Delta_0)}{k_B T} \gg 1$ and limiting to positive biases only we can solve the integrals of Eq. 3:

$$\begin{aligned} \int_{-\infty}^{+\infty} f(E, eV) [1 - f(E + \Delta E, eV)] dE &= \\ \Delta E + \frac{1}{4} \frac{\Delta E - eV}{e^{(\Delta E - eV)/k_B T} - 1}, & \text{when } \Delta E > 0, \\ \frac{1}{4} \frac{-\Delta E - eV}{e^{(-\Delta E - eV)/k_B T} - 1}, & \text{when } \Delta E < 0 \end{aligned} \quad (4)$$

With the results of Eq. (4) we can calculate the rates of Eq. (2). The set of linear equations can be solved analytically and the resulting occupation numbers can be inserted in the expression for the current (1). Figure 4(a) shows a measurement and fit for a H_2 molecule in a Pt contact, while Fig. 4(b) shows the evolution of the occupation numbers, which clearly shows a sharp drop in n_0 and rise in n_1 once the vibration mode energy is reached.

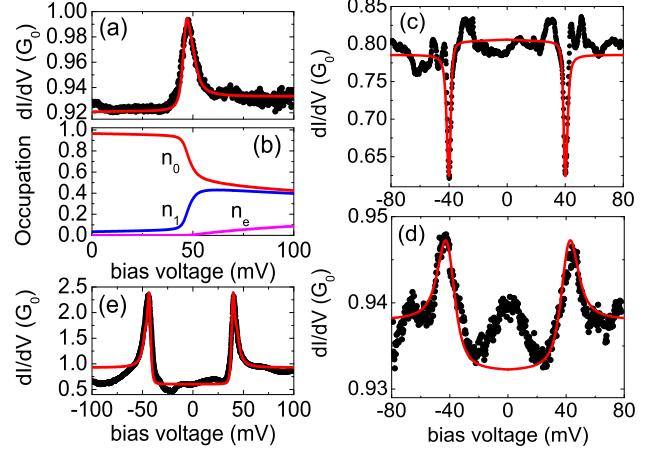


FIG. 4: (Color online) Several fits of VITLS for single molecule junction contacts and atomic chains: (a) H_2 in a Pt contact; (b) Evolution of occupation numbers n_0 (red), n_1 (blue), n_e (magenta) for the fit of (a); (c) H_2 in a Au atomic chain; (d) H_2 in Au atomic chain; (e) O_2 in a Ni contact. The experimental temperature was 5 K in all cases except in (d) where the contact was heated to 20 K. For fit parameters see [6].

Figure 4(c) and (d) show spectra and fits for H_2 in a Au atomic chain at different temperatures. In Fig. 4(d) the regular phonon mode of the gold atoms in the chain at ± 15 meV can be clearly resolved next to the VITLS spikes. Figure 4(c) shows clear negative spikes even when large conductance fluctuations are present. These conductance fluctuations arise from interference of electron waves that scatter on defects and impurities close to the contacts [16]. Often, the regular vibration mode signal is masked by these fluctuations for conductances $< 1 G_0$, but the spikes clearly show. Figure 4(e) shows a spectrum and fit for O_2 in a Ni contact, and a fit was also added in Fig. 1(a), indicating that the model is not limited to hydrogen. The energy splitting between the two ground states Ψ_0 and Ψ_1 is 1–3 meV and is clearly larger than the fitted temperatures, which are about 0.3–0.6 meV (~ 4 –7 K) in all cases except for the spectrum of Fig. 4(d), for which the fitted temperature is 1.6 meV (~ 20 K) in agreement with the bath temperature. This implies that for $V \ll V_m$ the system is almost entirely in state Ψ_0 (see Fig. 4(b)), which is a prerequisite for obtaining a sufficiently large decrease of n_0 and increase of n_1 when V reaches V_m .

The model described above involves vibrational excitation above the barrier. In the case of hydrogen it is also possible that the molecule tunnels through the reduced effective barrier after excitation. We have elaborated a similar model for the case involving a weak tunnel coupling between excited states in the potential wells, which leads to the emergence of two superpositions of the excited states [17]. The resulting fits are not qualitatively

different. Furthermore, we have investigated the possibility of double excitations in the potential wells within the framework of our model. The resulting fits could not reproduce the experimental data, due to a much slower population of the states of well 1.

As illustrated by Fig. 2 the conditions for obtaining a measurable VITLS signal can be obtained by control of the MCBJ: It is possible by pulling and pushing the contact very gently to explore the potential landscape for the molecule so that the proper conditions for a VITLS can be obtained.

Recently STM measurements on a hydrogen covered Cu(111) surface yielded very similar anomalous dI/dV spikes [18]. The mechanism is probably related but not identical to ours because the peak positions showed strong dependence on hydrogen coverage of the surface. Moreover, the observed two-level fluctuations were observed to remain in the σ_1 -state above the threshold voltage. Possibly the geometry of the experiment allows for the molecules to be excited into positions away from the STM tip, after which they are slow to return.

The IETS measurements by Wang *et al.* [19] on a self-assembled monolayer of octanedithiols have been discussed [20] in view of the unexpected shape of the features in d^2I/dV^2 . Several of the stronger features have the shape of peaks in the first derivative, dI/dV , but still at the energies of known vibration mode energies for the molecules. Our model produces a good fit to the data by Wang *et al.* but the temperature in the fit is much higher than the experimental temperature. This may be attributed to inhomogeneous broadening in the ensemble of molecules measured.

We conclude that the observed spikes in dI/dV measurements on single molecule junctions are most probably due to vibrationally induced two-level systems. Our simple model fits the data very well. It will be interesting to investigate potential energy landscapes for single molecule junctions in atomistic models in order to obtain a more concrete picture of the states involved. The VITLS acts as an intrinsic amplification mechanism providing a strong spectroscopic signature for local vibration modes and the anomalous spikes in dI/dV are expected to be a general feature of single molecule junctions. It can therefore be exploited as a new spectroscopic tool in analogy to the "action spectrum" proposed by Komeda *et. al.* [12]. Specifically, Au atomic chains with H_2 , that very rarely show regular vibration mode spectra, have a vibration mode around 42 meV, as seen from Fig. 2(d).

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